

Chaotic motion in the Johannsen-Psaltis spacetime

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- "No-hair" theorem
- Kerr hypothesis

The Kerr spacetime: symmetries \implies integrals of geodesic motion:

- stationarity \implies energy E ,
- axisymmetry \implies z-component of angular momentum L_z ,
- a "hidden symmetry" \implies Carter constant \mathcal{K} .

Thus, geodesic motion in the Kerr spacetime background is integrable.

Integrability broken by perturbation:

- increasing the number of degrees of freedom,
- reducing the system's degree of symmetry.

"Non-Kerr" spacetime:

- family of spacetimes which parametrically deviate from the Kerr spacetime,
- typical feature of non-Kerr spacetime: geodesic non-integrability.

Detection:

- gravitational wave signals,
- electromagnetic signals,

from EMRIs (Extreme Mass Ratio Inspiral).

Analysis methods:

- frequency analysis,
- recurrence analysis.

Geodesic motion in the Johannsen-Psaltis spacetime

- Metric:

$$ds^2 = g_{tt}dt^2 + g_{rr}dr^2 + g_{\theta\theta}d\theta^2 + g_{\phi\phi}d\phi^2 + 2g_{t\phi}dtd\phi \quad (1)$$

- Parameters: $M, a, \epsilon_2, \epsilon_3, \epsilon_4 \dots$
- Motion generated by the action

$$\mathcal{S} = \int_{\tau_1}^{\tau_2} \mathcal{L} d\tau, \quad \mathcal{L}(x^\mu, \dot{x}^\mu) = \frac{m}{2} g_{\mu\nu} \dot{x}^\mu \dot{x}^\nu \quad (2)$$

- Autonomous system (evolution parameter τ) \implies conservation of the Hamiltonian

$$\mathcal{H} = \frac{1}{2m} g^{\mu\nu} p_\mu p_\nu \quad (3)$$

- Integrals \mathcal{H}, E, L_z (and in the Kerr case \mathcal{K}) are independent and in involution ($\{I_i, I_j\} = 0$)

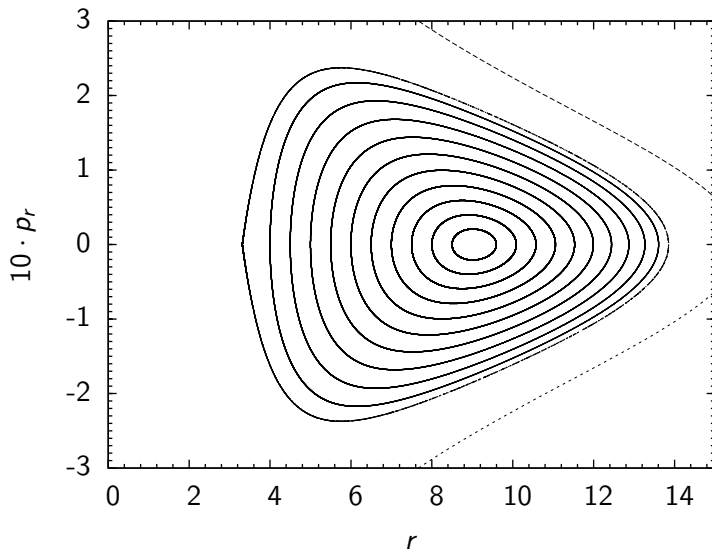
Calculated using:

- program written in C,
- RK4 algorithm,
- accuracy tracked by $\Delta\mathcal{L}_{rel}$.

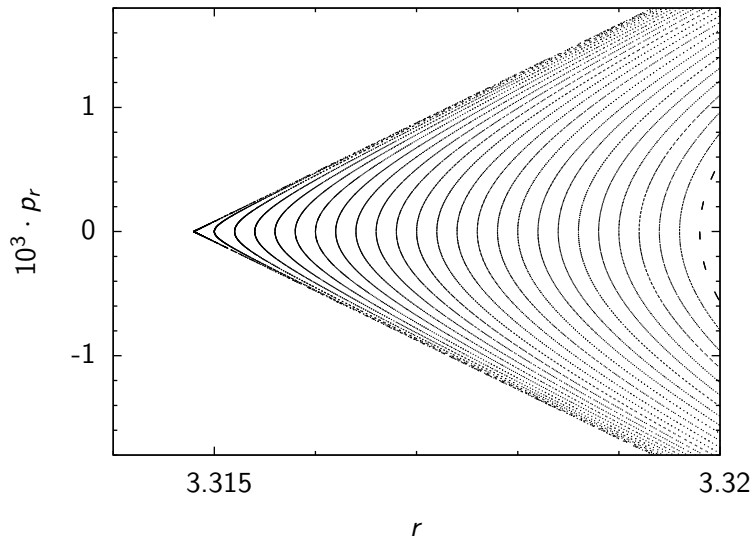
Surface of section

- A plane in the phase space of a system
- Used here: the equatorial plane $\theta = \pi/2$ with $\dot{\theta} \geq 0$
- Liouville-Arnol'd theorem: nested invariant circles in integrable case

Surface of section in an integrable system



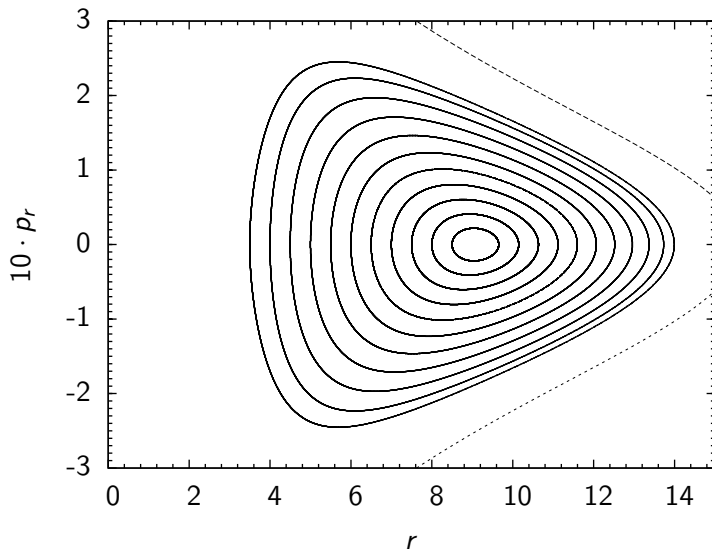
Surface of section in an integrable system



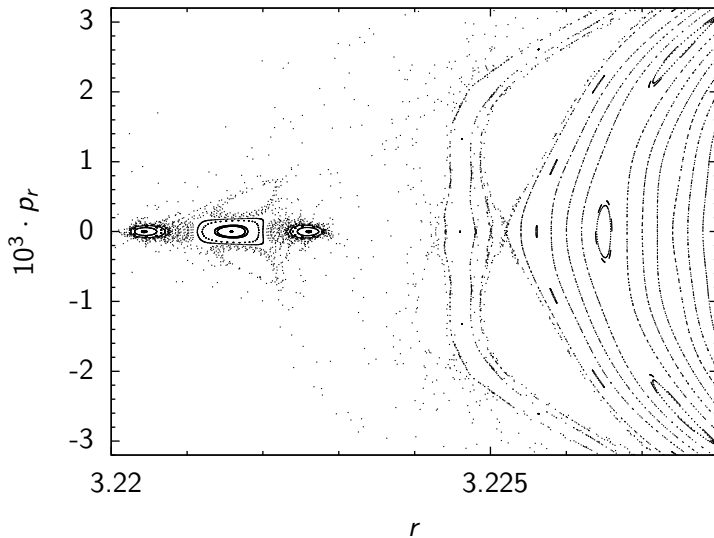
Surface of section in the perturbed system

- KAM theorem: some circles remain after perturbation
- Poincare-Birkhoff theorem: island chains

Surface of section in the perturbed system



Surface of section in the perturbed system



Rotation number

Regular trajectories:

- fundamental frequencies ω^r, ω^θ
- rotation number $\omega = \omega^r / \omega^\theta$

Take angles between consecutive points: $\vartheta_n = \text{ang}[x_n, x_{n+1}]$

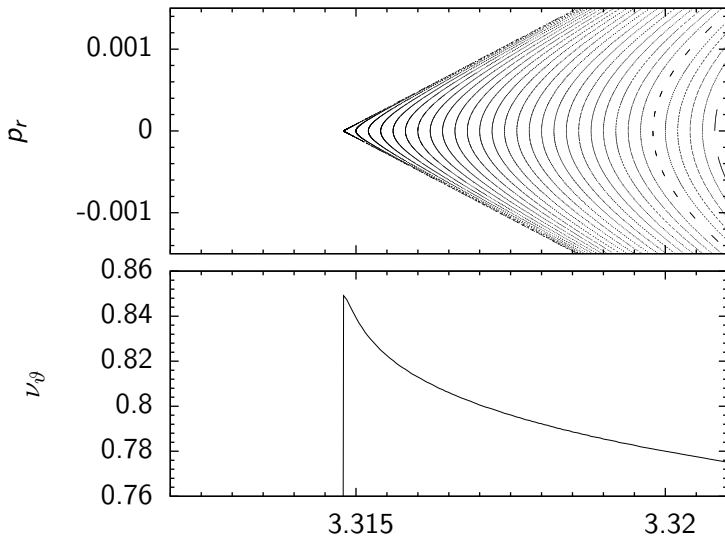
- for regular trajectories

$$\nu_\vartheta = \lim_{N \rightarrow \infty} \frac{1}{2\pi N} \sum_{i=1}^N \vartheta_i = \omega \quad (4)$$

- alternative: frequency analysis (Laskar, 1993)

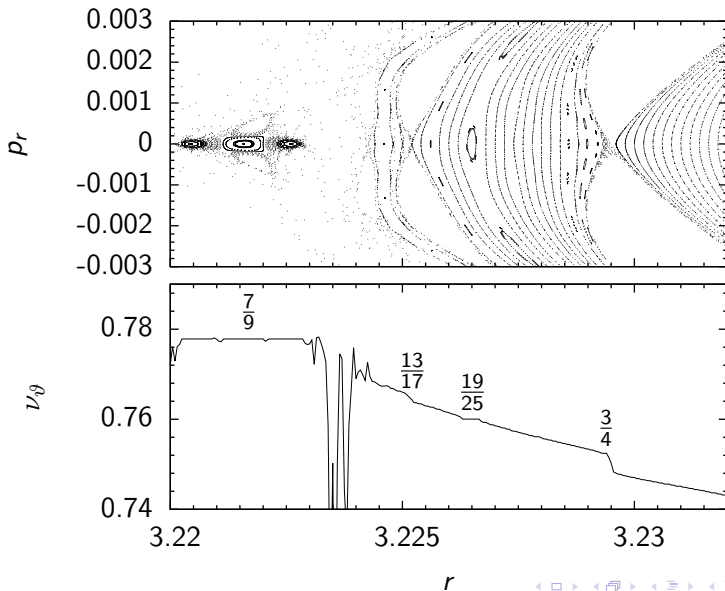
- ν_θ can be calculated for a chaotic trajectory as well
- in an integrable system: strictly monotonic function of the initial condition
- non-integrable system: non-monotonic variations when passing through chaotic zones

Rotation number in an integrable system



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Rotation number in the perturbed system



Deviation vector: between two infinitesimally close trajectories

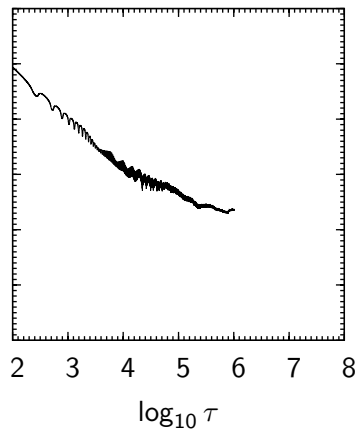
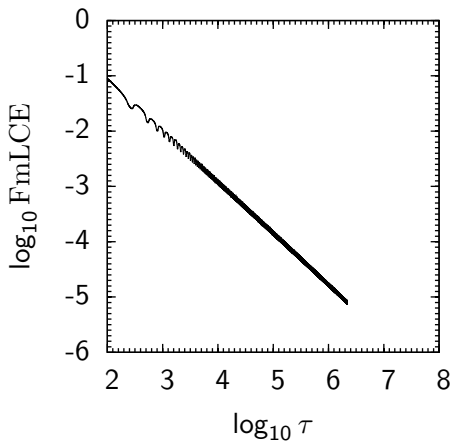
$$\ddot{\xi}^\mu + \frac{\partial \Gamma^\mu_{\kappa\lambda}}{\partial x^\nu} \dot{x}^\kappa \dot{x}^\lambda \xi^\nu + 2\Gamma^\mu_{\kappa\lambda} \dot{x}^\kappa \dot{\xi}^\lambda = 0 \quad (5)$$

$$\Xi^2 = g_{\mu\nu} \xi^\mu \xi^\nu \quad (6)$$

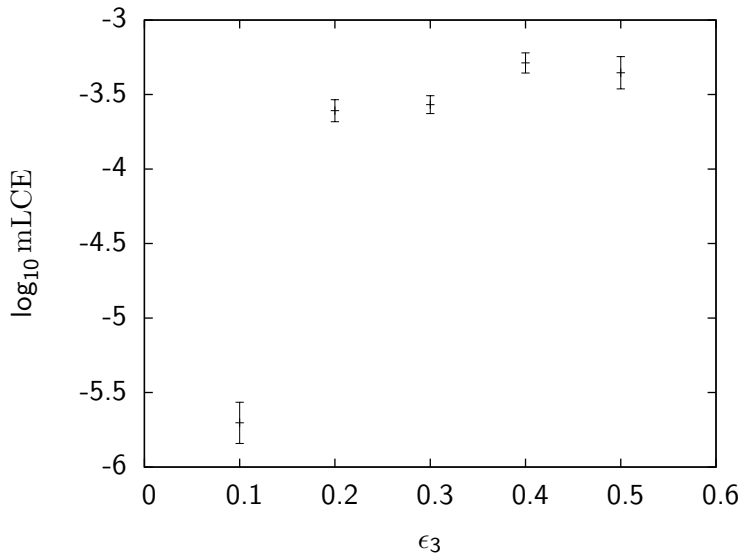
$$\text{mLCE} = \lim_{\tau \rightarrow \infty} \frac{1}{\tau} \log \left[\frac{\Xi(\tau)}{\Xi(0)} \right] \quad (7)$$

- regular trajectory: $\Xi(\tau) \sim \tau$, $\text{mLCE} = 0$
- chaotic trajectory: $\Xi(\tau) \sim e^{\lambda\tau}$, $\text{mLCE} = \lambda$

Lyapunov exponents



Lyapunov exponents



It was shown by clear numerical examples using different methods that geodesic motion in the Johannsen-Psaltis spacetime corresponds in general to a non-integrable system.

Time for discussion

Any questions?