

**At Centenary of Einstein's  
General Theory of Relativity:  
Testing the Nature of Accreting  
Black Holes**

*Selected topics as science drivers for mission proposals*

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# Objects and models

- Active galactic nuclei
- Stellar-mass black holes
- Intermediate-mass black holes (?)

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# Objects and models

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  - ...time-dependent, non-axisymmetric
- GR effects taken into account
- Link to a spectrum-fitting procedure → parameters

# E1.-vac. soln. w/rotating BH

*Exact solution: magnetized Kerr-Newman BH:*

$$g = |\Lambda|^2 \Sigma (\Delta^{-1} dr^2 + d\theta^2 - \Delta A^{-1} dt^2) \\ + |\Lambda|^{-2} \Sigma^{-1} A \sin^2 \theta (d\phi - \omega dt)^2,$$

$\Sigma = r^2 + a^2 \cos^2 \theta$ ,  $\Delta = r^2 - 2Mr + a^2 + e^2$ ,  
 $A = (r^2 + a^2)^2 - \Delta a^2 \sin^2 \theta$  are functions from the  
Kerr-Newman metric.

$\Lambda = 1 + \beta\Phi - \frac{1}{4}\beta^2\mathcal{E}$  is given in terms of the Ernst complex  
potentials  $\Phi(r, \theta)$  and  $\mathcal{E}(r, \theta)$ :

$$\Sigma\Phi = ear \sin^2 \theta - \Im e (r^2 + a^2) \cos \theta, \\ \Sigma\mathcal{E} = -A \sin^2 \theta - e^2 (a^2 + r^2 \cos^2 \theta) \\ + 2\Im a [\Sigma (3 - \cos^2 \theta) + a^2 \sin^4 \theta - re^2 \sin^2 \theta] \cos \theta.$$

Bretón Baez N., García Díaz A., J. Math. Phys. 27, 562 (1986)

# Thank you!

See Additional Material for Discussion:

High-frequency elmg. waves

Polarization tensor

Stokes parameters

# High-frequency elmg. waves

Basic equations – vacuum case:  $F^{\mu\nu}{}_{;\nu} = 0$ ,  $*F^{\mu\nu}{}_{;\nu} = 0$ .

$$E^\alpha = F^{\alpha\beta} u_\beta, \quad *F_{\mu\nu} \equiv \frac{1}{2} \varepsilon_{\mu\nu}{}^{\rho\sigma} F_{\rho\sigma}$$

An electromagnetic wave is an approximate test-field solution of the Maxwell equations:

$$F_{\alpha\beta} = \Re e [u_{\alpha\beta} e^{\mathfrak{I}S(x)}].$$

A fixed background geometry is assumed.

- Phase  $S(x)$  ... rapidly varying function
- Amplitude  $u_{\alpha\beta}$  ... slowly varying function
- Wave vector  $k_\alpha \equiv S_{,\alpha}$  ... parallel transport, null geodesics

$$k_{\alpha;\beta} k^\beta = 0, \quad k_\alpha k^\alpha = 0.$$

# Polarization tensor

- Polarization tensor ...  $J_{\alpha\beta\gamma\delta} \equiv \frac{1}{2} \langle F_{\alpha\beta} \bar{F}_{\gamma\delta} \rangle$
- In observer's rest-frame ...  $J_{\alpha\beta} \equiv J_{\alpha\beta\gamma\delta} u^\beta u^\delta = \langle E_\alpha \bar{E}_\beta \rangle$
- Four parameters ...  $S_A \equiv \frac{1}{2} (k_\alpha u^\alpha)^2 F_A \quad (A = 0, \dots, 3)$   
( $F_A$  ... constructed by projecting onto a tetrad  $e_{(i)}^\alpha$ )

*“On the composition and resolution of streams  
of polarized light from different sources”*



- References: [1] Sir George Stokes (1852), *Trans. Cambridge Phil. Soc.*, 9, 399  
[2] Chandrasekhar (1950), *Radiative Transfer* (Oxford: Clarendon)  
[3] Cocke & Holm (1972), *Nature*, 240, 161  
[4] Jauch & Rohrlich (1955), *The Theory of Photons and Electrons* (Reading: Wesley)



# Stokes parameters

$$S_0 \equiv J_{\alpha\beta} \left( e_{(1)}^\alpha e_{(1)}^\beta + e_{(2)}^\alpha e_{(2)}^\beta \right) = \langle |E_{(1)}|^2 + |E_{(2)}|^2 \rangle$$

$$S_1 \equiv J_{\alpha\beta} \left( e_{(1)}^\alpha e_{(1)}^\beta - e_{(2)}^\alpha e_{(2)}^\beta \right) = \langle |E_{(1)}|^2 - |E_{(2)}|^2 \rangle$$

$$S_2 \equiv J_{\alpha\beta} \left( e_{(1)}^\alpha e_{(2)}^\beta + e_{(2)}^\alpha e_{(1)}^\beta \right) = \langle E_{(1)} \bar{E}_{(2)} + E_{(2)} \bar{E}_{(1)} \rangle$$

$$S_3 \equiv \Im J_{\alpha\beta} \left( e_{(1)}^\alpha e_{(2)}^\beta - e_{(2)}^\alpha e_{(1)}^\beta \right) = \Im \langle E_{(1)} \bar{E}_{(2)} - E_{(2)} \bar{E}_{(1)} \rangle$$

$S_1$ ,  $S_2$ , and  $S_3$  determine the polarization state.

References: [5] Anile (1989), *Relativistic fluids and magneto-fluids* (Cambridge)

[6] Madore (1974), *Comm. Math. Phys.*, 38, 103

[7] Bičák & Hadrava (1975), *A&A*, 44, 389

[8] Breuer & Ehlers (1980), *Proc. Roy. Soc. Lond. A*, 370, 389

[9] Broderick & Blandford (2003), *MNRAS*, 342, 1280